

# **ENERGY/COLOR FLOW**

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# OUTLINE

- Motivation
- Rapidity gap process
- Factorization and refactorization
- Color mixing matrices
- Resummation
- Gap fraction
- Summary

# MOTIVATION

- The role of perturbative QCD
  - Tevatron
    - QCD was a theory to be tested.
      - Discovery of top quarks
  - LHC
    - Precision
    - A tool to search for physics BSM

# MOTIVATION

- The role of perturbative QCD
  - Tevatron
    - QCD was a theory to be tested
      - Discovery of top quarks
  - LHC
    - Precision
    - A tool to search for physics BSM
    - A tool to extract properties in BSM

# MOTIVATION – BASIC PROPERTIES IN BSM

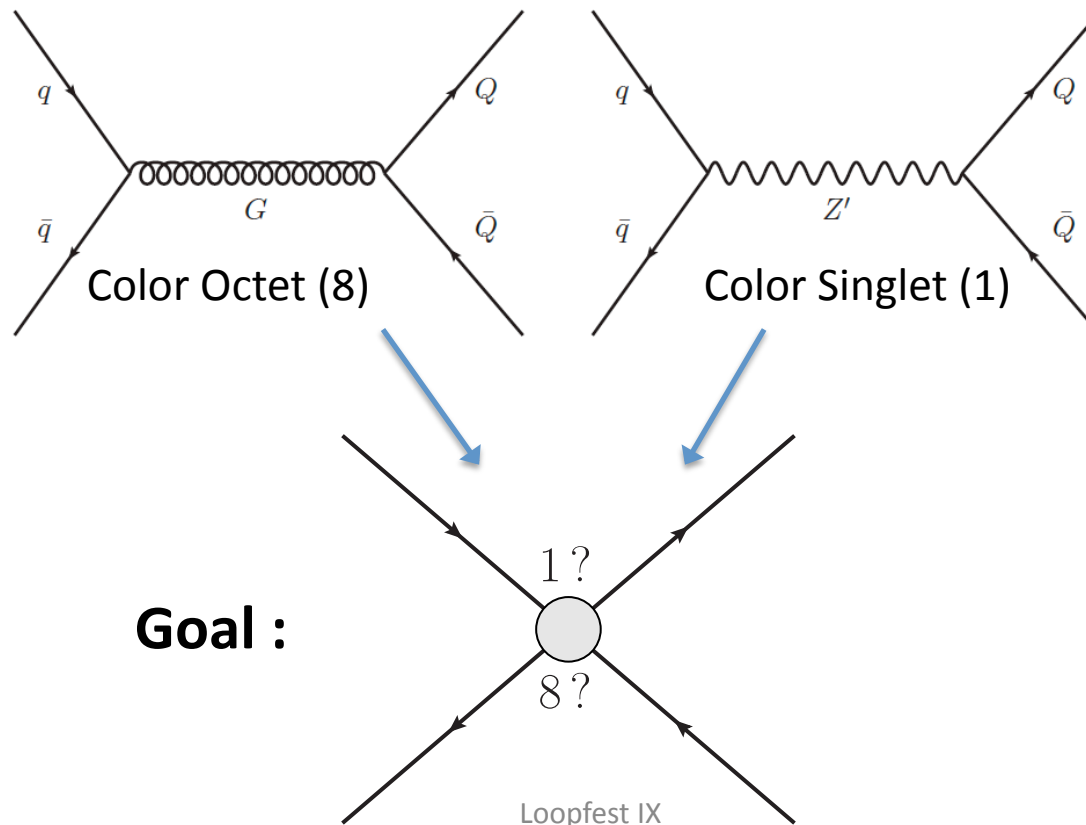
- There are many proposals for the extension of the SM.
  - The most common feature of the models is that they include new heavy particles.
  - They decay into the SM particles.
    - KK gauge boson  $\rightarrow$  a quark or lepton pair
      - [Arkani-Hamed, Dimopoulos, Dvali, 1998], [Randall, Sundrum, 1999]
    - $Z' \rightarrow HZ$
- The detailed analysis of products of the resonance decay allows one to study the properties in new theory.
- Yes, but how?

# MOTIVATION – SPECIFIC EXAMPLES WITH FOLLOWING CONSIDERATIONS

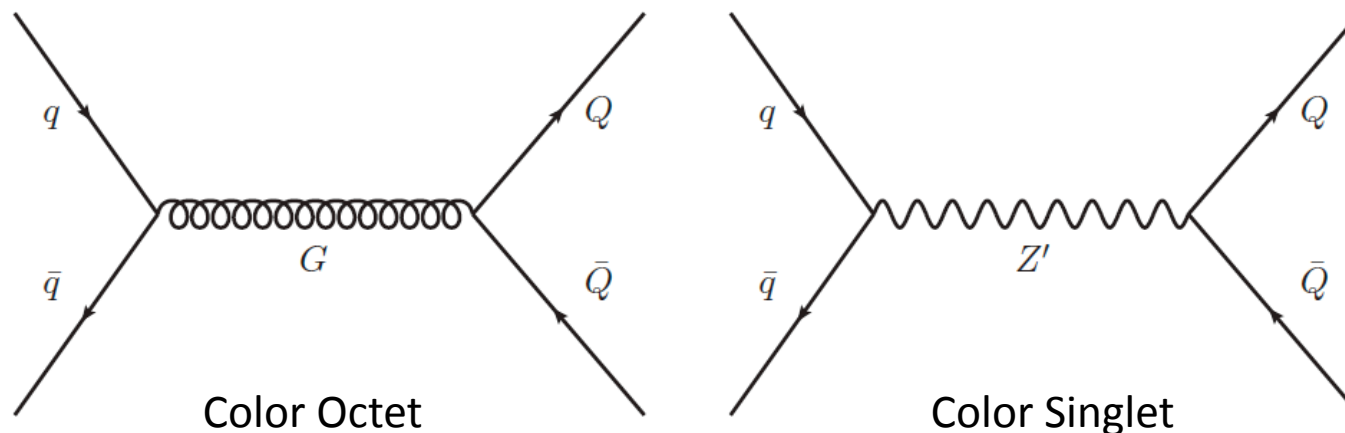
- 1) In many scenarios, heavy particles are created with  $O(\text{TeV})$  mass, and decay into the SM particles.
- 2) The resonances are coupled strongly to heavy partons, compared to their couplings to light partons and leptons.
- 3) Let's assume that we succeed in distinguishing resonances from background events
  - Invariant mass distributions and many other techniques

- What is next?
  - The SM gauge content of resonance particles
- Based on the previous considerations, we guess two possible resonance processes.

- What is the next?
  - The SM gauge content of resonance particles
- Based on the previous considerations, we guess two possible resonance processes. (s-channel due to strong couplings to heavy quarks)

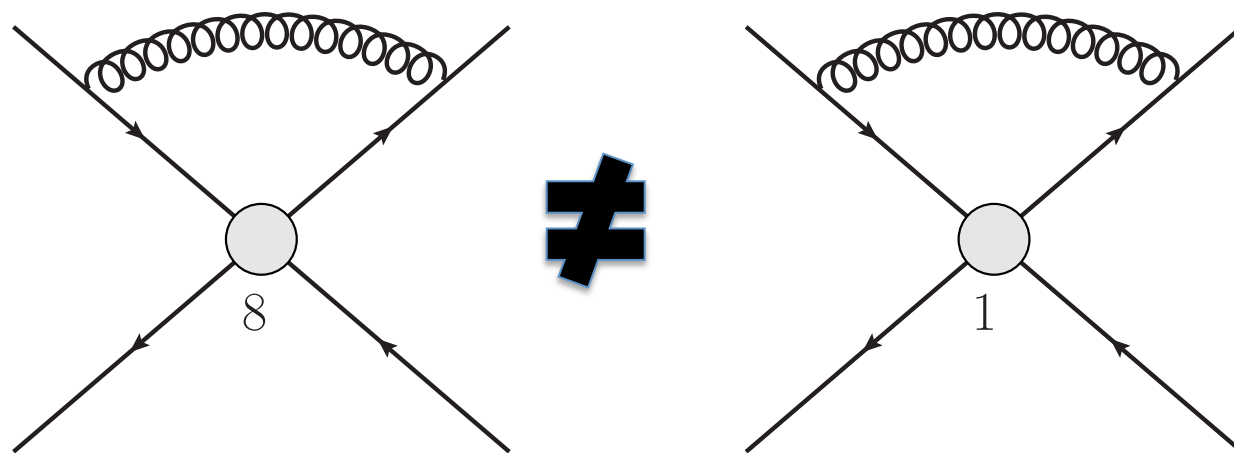






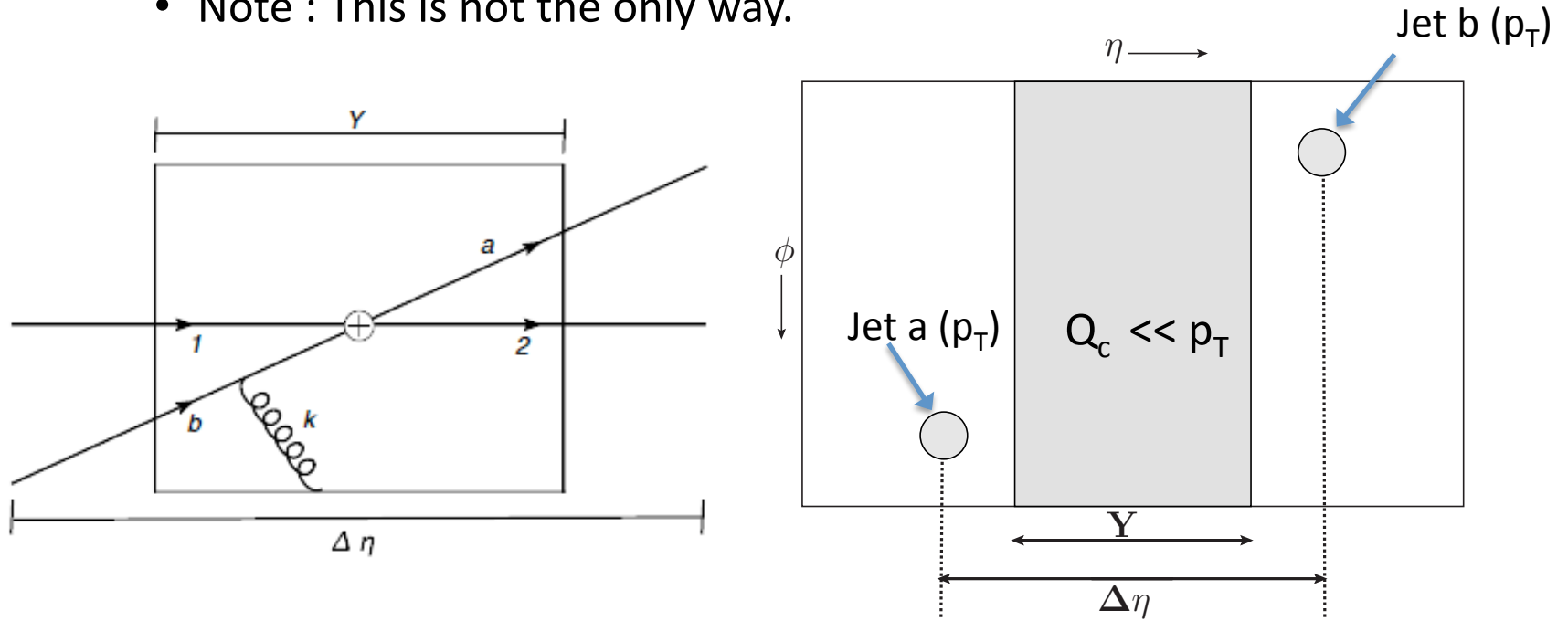
- Observation
  - All external particles are SM SU(3) particles.
  - Color can flow only through octet resonances ( $G$ ).
- What is the physical quantity distinguishing one from another? And how can we calculate it?

- The difference induced by different color flow in processes begins to appear at higher orders
  - Ex. We measure energy ( $Q_c$ ) of gluon emissions into a certain area of a detector



- The different patterns of radiation may appear, depending on the  $SU(3)$  gauge content of resonances.

- How can we quantify the amount of radiation?
  - Rapidity gap processes [Oderda, Sterman (1998)], [Dasgupta, Salam (2001)], [Appleby, Seymour (2003)]
    - Note : This is not the only way.



- We require to tag inclusively a pair of jets or heavy quarks and measure energy in a systematic central region, spanning rapidity  $Y$ .

- Our candidate quantity - inclusive top pair production cross-sections with energy in the central region

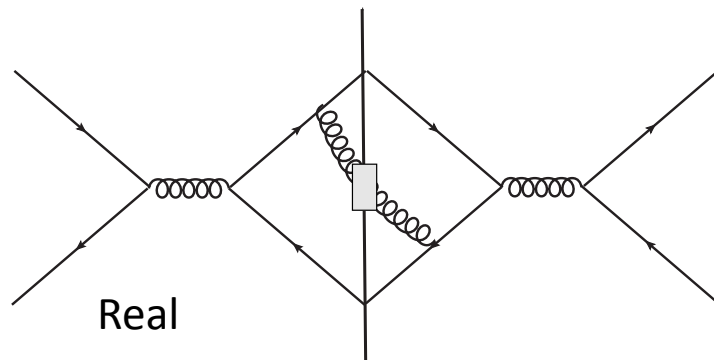
- equal to  $Q_c$   $\frac{d\hat{\sigma}^{(f)}}{d\Delta\eta dQ_c}(M, m_t, Q_c, \mu_F, \Delta\eta, Y, \alpha_s(\mu_F))$

- or up to energy flow  $Q_0$

$$\frac{d\hat{\sigma}^{(f)}}{d\Delta\eta}(M, m_t, Q_0, \mu_F, \Delta\eta, Y, \alpha_s(\mu_F)) = \int_0^{Q_0} \frac{d\hat{\sigma}^{(f)}}{d\Delta\eta dQ_c}(M, m_t, Q_c, \mu_F, \Delta\eta, Y, \alpha_s(\mu_F))$$

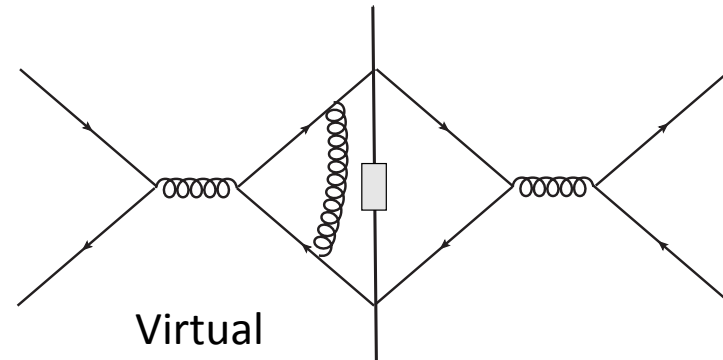
- When the latter is divided by total cross section, the ratio can be interpreted as a probability of a top pair production with radiation into the gap up to energy flow  $Q_0$ .

- Compare these probabilities, called “gap fractions”, for color singlet and octet resonances.
- How can we calculate a gap fraction ?
  - Partonic cross section for each resonance process at fixed order.
  - MC simulations



Real

$$= \int_{\Omega_{\text{gap}}} d\text{PS } \omega_R$$



Virtual

$$= \int_{\Omega} d\text{PS } \omega_V$$

$$\int_0^{Q_0} \frac{dk_0}{k_0} \int d\Omega_{\text{gap}} (c_1 \omega_R) - \int_0^{\infty \rightarrow \mu} \frac{dk_0}{k_0} \int d\Omega (c_2 \omega_V) = C(Y) \ln \left( \frac{\mu}{Q_0} \right)$$

- However, we can do better than this by using factorization and evolution equations
  - Each loop produces a single log due to soft gluon exchange. [PRD17:2773,2789, Sterman (1978)]
  - Resummation; the sum of all leading enhancements along with associated color flow to all orders
- We measure observable from hadronic collisions. Thus we need a connection between partonic and hadronic cross sections.
  - Factorization [Adv.Ser.Dir.HEP, Collins, Soper and Sterman (1985)]

$$\frac{d\sigma_{AB}}{d\Delta\eta dQ_c} = \sum_{f_1, f_2} \int dx_1 dx_2 \phi_{f_1/A}(x_1, \mu_F) \phi_{f_2/B}(x_2, \mu_F) \frac{d\hat{\sigma}^{(f)}}{d\Delta\eta dQ_c}$$

# REFACTORIZATION

- In general, we can perform a further factorization on the partonic rapidity gap cross section

$$\frac{d\hat{\sigma}^{(\text{f})}}{d\Delta\eta}(M, m_Q, Q_0, \mu_F, \Delta\eta, Y, \alpha_s(\mu_F)) = \sum_{L,I} H_{IL}^{(\text{f})}(M, m_Q, \mu_F, \mu, \Delta\eta, \alpha_s(\mu)) \times S_{LI}^{(\text{f})}\left(\Delta\eta, Y, \frac{Q_0}{\mu}, \alpha_s(\mu), m_Q\right)$$

- H represents physics at the large momentum scale
  - the mechanism of a heavy particle resonance is contained in H
- S describes the infrared behavior of color exchange starting at short distance, and represents the radiation of soft gluons up to the scale  $Q_0$ .
- I or L indicates some basis of color tensors linking external partons.

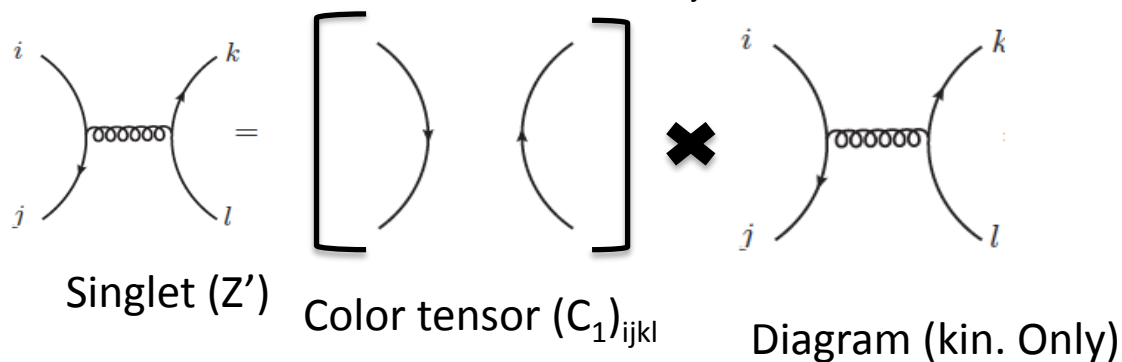
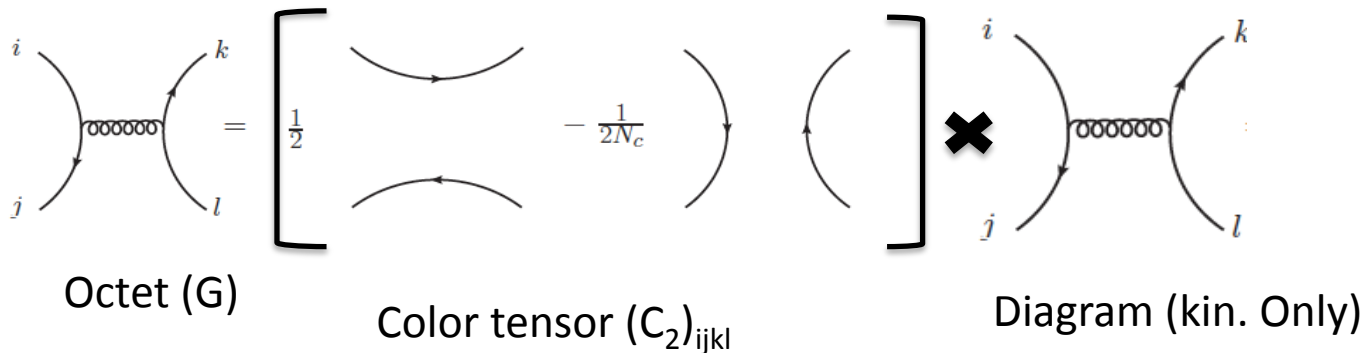
# COLOR MIXING MATRIX

- How can we describe color flow ?
  - In fixed order calculations, we take amplitudes squared.
    - It is difficult to track color flow.
    - It is nontrivial to imagine color flow at higher orders.
  - Our goal : Describe color flow in the space of color tensors under exchange gluons  
[Sen (1983),Botts, Sterman (1989), Kidonakis, Orderda , Sterman(1998)]
    - Note: this is a key for resummation



- Amplitudes for our processes are four index tensors.
- Let's decouple color tensors from diagrams.
  - Use the color identity

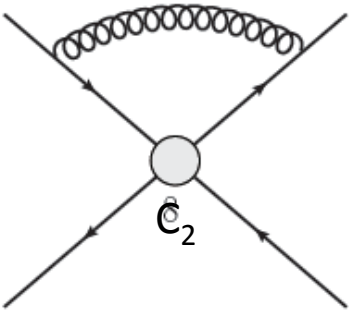
$$(T_F^a)_{ji}(T_F^a)_{kl} = \frac{1}{2} \left( \delta_{ki}\delta_{jl} - \frac{1}{N_c} \delta_{ji}\delta_{kl} \right)$$



**Resonance amplitude in color spaces -> hard function**

# COLOR MIXING MATRIX

- Let's study how exchange of a gluon mixes the color tensors
  - If color tensor basis  $c_1$  and  $c_2$  span the entire color space, the result of color exchanges would be a linear function of two basis tensors.



$$= \left[ \frac{1}{2} \begin{array}{c} \text{gluon exchange} \\ \text{diagram} \end{array} - \frac{1}{2N_c} \begin{array}{c} \text{gluon exchange} \\ \text{diagram} \end{array} \right] \times \begin{array}{c} \text{gluon exchange} \\ \text{diagram} \end{array}$$

$$= \left[ -C_F/(2N_c)(C_1)_{ijkl} - (C_F - 1/2N_c^2)(C_2)_{ijkl} \right] \times \begin{array}{c} \text{gluon exchange} \\ \text{diagram} \end{array}$$

# COLOR MIXING MATRIX

- At one loop order, the color mixing matrix is given by [Berger, Kucs, Sterman (2001)]

$$C = \begin{pmatrix} C_F \left( \text{diagram 1} + \text{diagram 2} \right) & \frac{C_F}{2N_c} \left( \text{diagram 3} + \text{diagram 4} + \text{diagram 5} + \text{diagram 6} \right) \\ \left( \text{diagram 7} + \text{diagram 8} + \text{diagram 9} + \text{diagram 10} \right) & C_F \left( \text{diagram 11} + \text{diagram 12} \right) - \frac{1}{2N_c} \left( 2 \left( \text{diagram 13} + \text{diagram 14} \right) + \text{diagram 15} + \text{diagram 16} \right) \end{pmatrix}$$

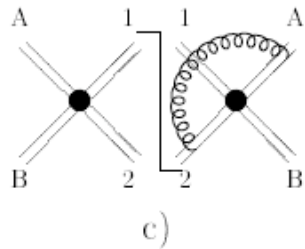
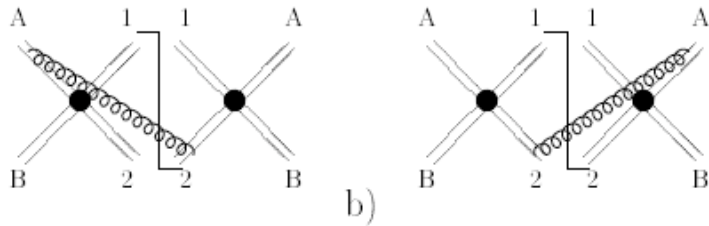
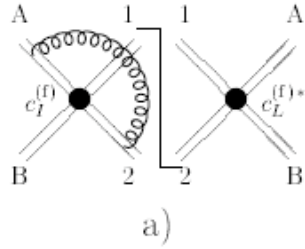
# RESUMMATION

- Now let's try to sum leading logarithms at each order to all orders [Oderda, Sterman (1998), Berger, Kucs, Sterman (2001)]

$$\frac{d\hat{\sigma}^{(\text{f})}}{d\Delta\eta}(M, m_Q, Q_0, \mu_F, \Delta\eta, Y, \alpha_s(\mu_F)) = \sum_{L,I} H_{IL}^{(\text{f})}(M, m_Q, \mu_F, \mu, \Delta\eta, \alpha_s(\mu)) \times S_{LI}^{(\text{f})}\left(\Delta\eta, Y, \frac{Q_0}{\mu}, \alpha_s(\mu), m_Q\right)$$

- Independence of the refactorization scale  $\mu$   
 -> Evolution equations in  $\mu$

$$\mu \frac{d}{d\mu} S_{LI} = -(\Gamma_S)_{LJ}^\dagger S_{JI} - S_{LI} (\Gamma_S)_{JI}$$



- Eikonal calculation
  - Effectively captures only divergences
  - Scaleless integrals (UV+IR=0)

$$\gamma_S^{(1)}(ij) = \frac{\omega_{(ij)}(-2\varepsilon)}{(\alpha_s/\pi)}$$

$$\omega_{(ij)} = \omega_V(ij) + \omega_R(ij)$$

$$\omega_{(ij)} = \left( \frac{\alpha_s}{\pi} \right) \left( -\delta_i \delta_j \Delta_i \Delta_j \frac{1}{2} \frac{1}{2\varepsilon} \int \frac{dy d\phi}{2\pi} \Theta(\vec{k}) \Omega_{ij} + \delta_i \delta_j \Delta_i \Delta_j \frac{i\pi}{2\varepsilon} \frac{(1 - \delta_i \delta_j)}{2} \right)$$

$$\Omega_{ij} = \frac{(p_i \cdot p_j) k_T^2}{(p_i \cdot k)(p_j \cdot k)}$$

# RESUMMATION

- Diagonalize color bases :

$$\begin{aligned} (\Gamma_S^{(f)}(\Delta\eta, Y, \rho))_{\gamma\beta} &\equiv \lambda_\beta^{(f)}(\Delta\eta, Y, \rho) \delta_{\gamma\beta} \\ &= (R^{(f)})_{\gamma I} (\Gamma_S^{(f)}(\Delta\eta, Y, \rho))_{IJ} (R^{(f)})^{-1}_{J\beta} \end{aligned}$$

- Evolution equation becomes

$$\mu \frac{d}{d\mu} S_{\gamma\beta} = -(\lambda_\gamma^* + \lambda_\beta) S_{\gamma\beta}$$

- Solution to this equation is

$$S_{\gamma\beta}^{(f)}\left(\Delta\eta, Y, \frac{Q_0}{\mu}, \alpha_s(\mu), \rho\right) = S_{\gamma\beta}^{(f)}(\Delta\eta, Y, 1, \alpha_s(Q_0), \rho) \exp\left[-E_{\gamma\beta}^{(f)} \int_{Q_0}^{\mu} \frac{d\mu'}{\mu'} \left(\frac{\beta_0}{2\pi} \alpha_s(\mu')\right)\right]$$

$$E_{\gamma\beta}^{(f)}(\Delta\eta, Y, \rho) = \frac{2}{\beta_0} [\lambda_\gamma^{(f,1)*}(\Delta\eta, Y, \rho) + \lambda_\beta^{(f,1)}(\Delta\eta, Y, \rho)]$$

- Could also do numerically
  - [Banfi, Salam, Zanderighi (2002), Almeida, Sterman, Vogelsang (2008)]

# RESUMMATION

- The cross section for heavy quark production with soft gluon emission of energy up to  $Q_0$  becomes

$$\frac{d\hat{\sigma}_{G/Z'}^{(f)}}{d\Delta\eta} = \sum_{\beta,\gamma} (H_{G/Z'}^{(f,LO)})_{\beta\gamma}(M, m_Q, \Delta\eta, \alpha_s(p_T)) S_{\gamma\beta}^{(f,0)} \left[ \frac{\ln\left(\frac{Q_0}{\Lambda}\right)}{\ln\left(\frac{p_T}{\Lambda}\right)} \right]^{E_{\gamma\beta}^{(f)}}$$

$M$  : Resonance mass

$Q_0$  : Energy of radiation into the gap

- Model dependence is encapsulated in a hard function.
- $H$  and  $S$  are connected only through color indices.

# GAP FRACTION

- The ratio of the number of events for heavy quark production with a specified rapidity gap to the total number of pair production events

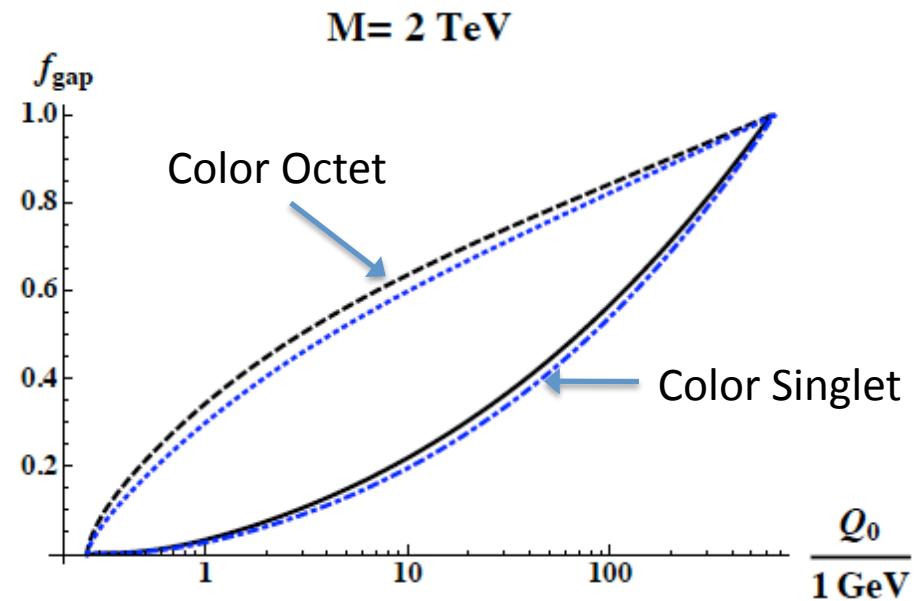
$$f_{gap}^{(LO)}(M, Q_0) = \frac{\frac{d\hat{\sigma}^{(f)}(M, Q_0)}{d\Delta\eta}}{\frac{d\hat{\sigma}^{(f, LO)}(M, M)}{d\Delta\eta}}$$

$$\frac{d\hat{\sigma}^{(f, LO)}}{d\Delta\eta} = \sum_{\beta, \gamma} H_{\beta\gamma}^{(f, LO)}(M, m_Q, \Delta\eta, \alpha_s(p_T)) S_{\gamma\beta}^{(f, 0)}$$

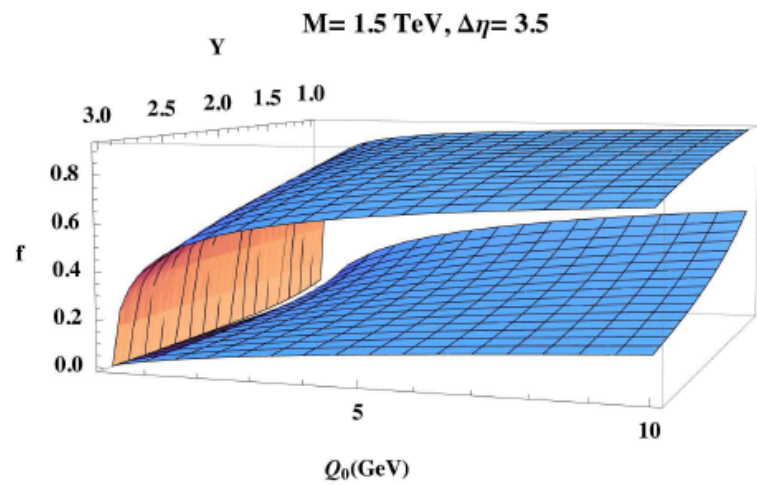
We don't need to know about the strength of couplings of new sectors to a quark pair !



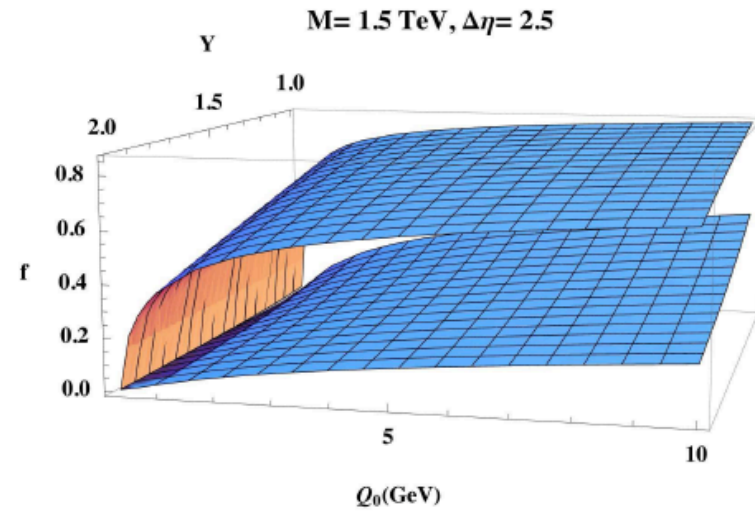
# GAP FRACTION



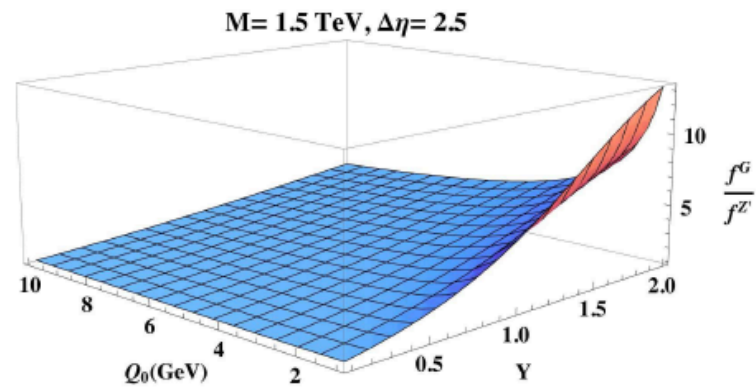
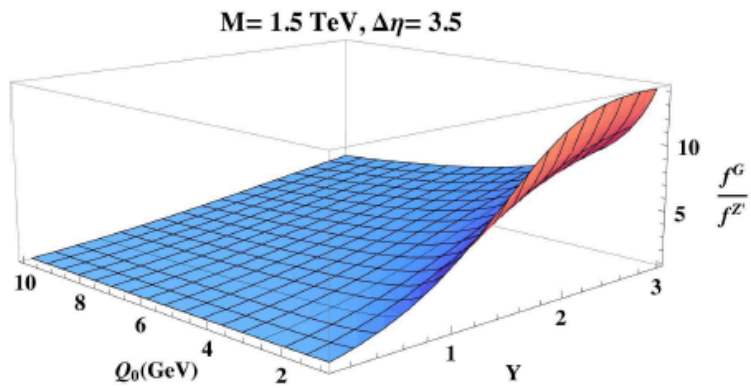
- More radiation for singlet than for octet resonances into the gap.



(a)



(b)



# SUMMARY

- Color flow is important information for BSM search
- We have shown that it is possible to determine the color  $SU(3)$  representation of resonance from NP
- More radiation for singlet than for octet resonances
- We can use any fixed region of rapidity and azimuthal angle rather than the simple rapidity gap
- We can use color flow information to distinguish signals from backgrounds at the LHC
  - Higgs or Color singlet object  $\rightarrow$  jets vs QCD jets

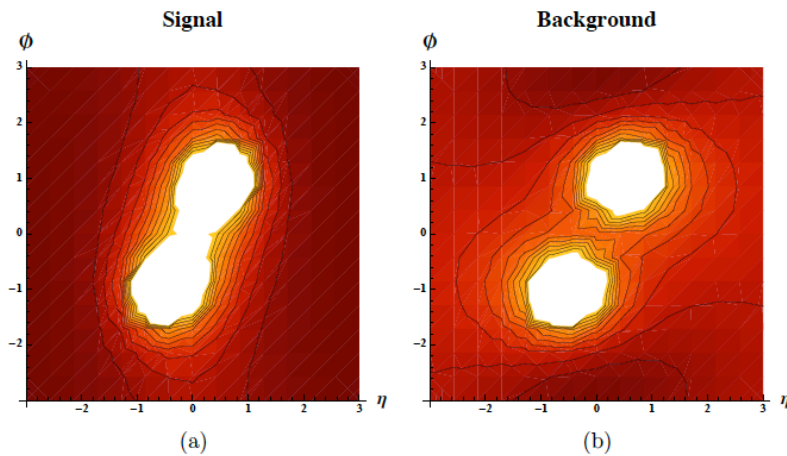
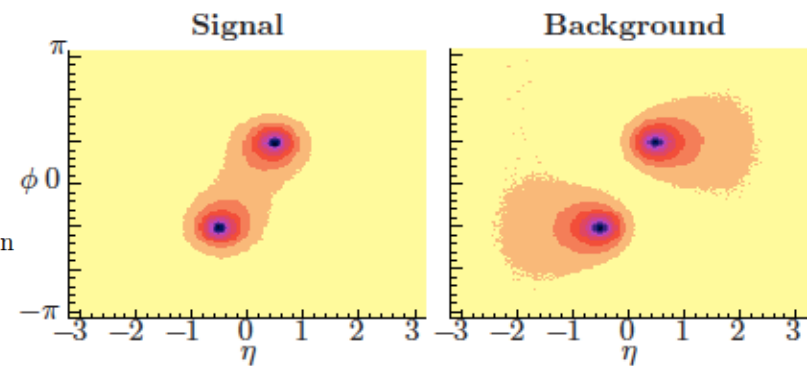


Figure 1: Density plots for  $\frac{d\sigma}{dQ_c d\eta d\phi}$  from Higgs vs from background process with gluon of  $Q_c$  into  $(\eta, \phi)$ . Brighter color, more probability to emit gluon into the area.

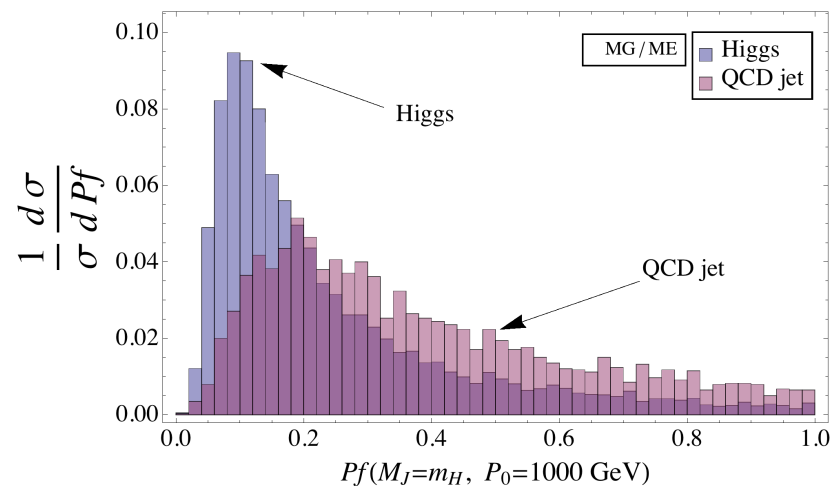
From analytical LL resummed result  
(left - work in progress,  $pp \rightarrow HZ \rightarrow b \bar{b} Z$ )

Comparing to the result from MC (below)  
[Gallicchio, Schwartz (2010)]

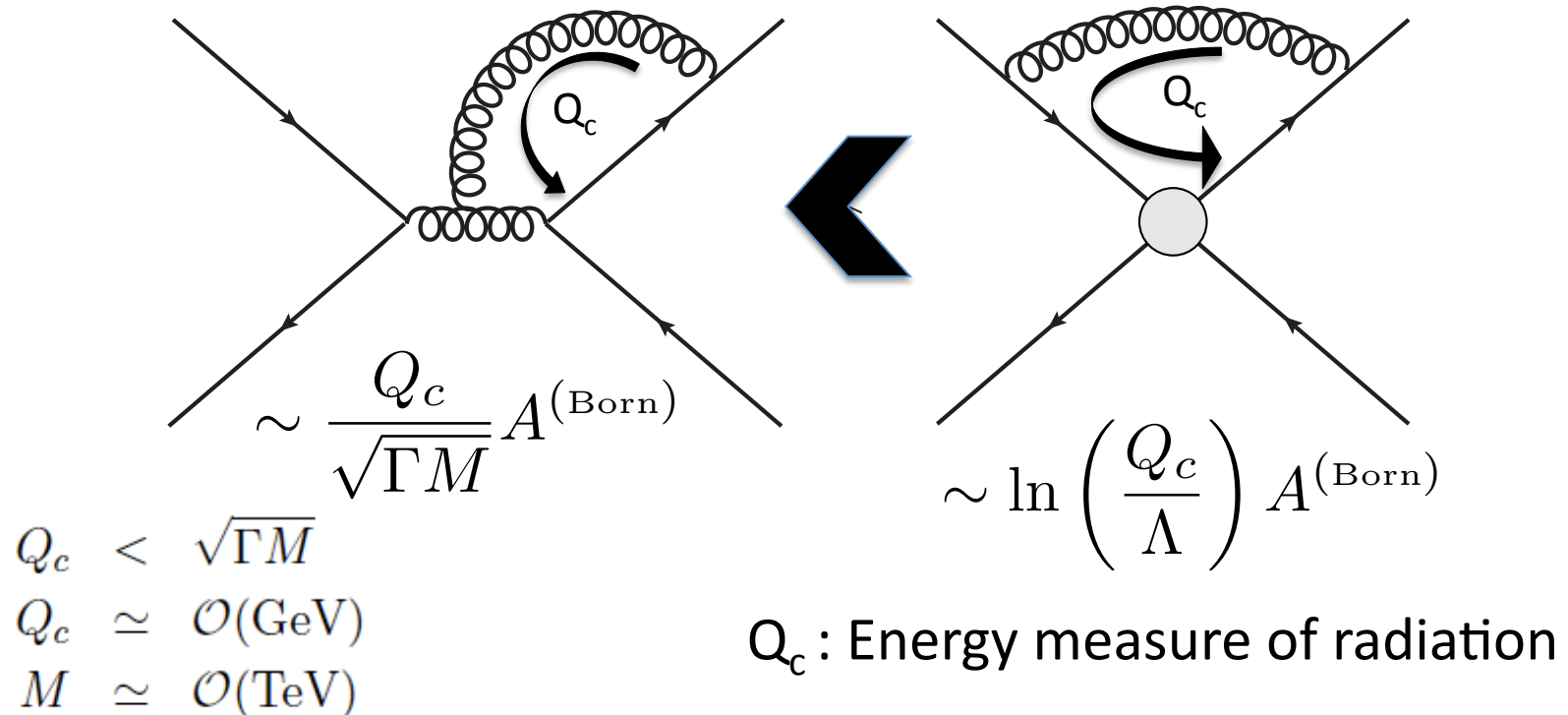


From MC simulations for  
Planar flow,

Almeida, Lee, Perez, Sterman, IS,  
Arxiv:1006.2035



- In this work, we consider resonance widths  $\Gamma > \mathcal{O}(\text{GeV})$  for TeV resonance masses.



$$\Gamma \simeq \frac{M}{6} \text{ for KK gluon in the RS models satisfying EW precision test [PRD, Agashe, et al (2008)]}$$

$$\Gamma \simeq \frac{M}{10} \text{ for axigluons and universal colorons [PRD, Choudhury, et al (2007)]}$$